

# Beyond-Hartree-Fock methods in CP2K: MP2, RPA and GW

## A tutorial

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- 1 Beyond-Hartree-Fock methods: Definition, benefits and limitations
- 2 The resolution of the identity (RI)
- 3 RI-MP2, RI-SOS-MP2 and RI-RPA for the correlation energy
- 4 Execution time of RI-MP2, RI-SOS-MP2 and RI-RPA
- 5 Inputs for RI-MP2, RI-SOS-MP2 and RI-RPA
- 6 Application: Density of water from RI-MP2 and RI-RPA Monte Carlo simulations
- 7 RI-GW for single-particle energies: Definition, input and application
- 8 Concluding remarks

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# The exchange-correlation term in Beyond-Hartree-Fock methods

Total energy  $E_{\text{tot}}$  of electrons in a molecule or periodic system:

$$E_{\text{tot}} = E_{\text{kin}} + E_{\text{el-core}} + E_{\text{Hartree}} + E_{\text{exchange (x)}} + E_{\text{correlation (c)}}$$

- **DFT:** exchange and correlation is a functional of the density  $n(\mathbf{r}) = \sum_{i=1}^{N_{\text{occ}}} |\psi_i(\mathbf{r})|^2$ :

$$E_{\text{xc}}^{\text{DFT}} = E_{\text{xc}}^{\text{DFT}}[n] \quad (\text{cheap})$$

- **Hartree-Fock (HF):** No correlation, exchange energy only depends on occupied MOs  $\psi_i(\mathbf{r})$ :

$$E_{\text{xc}}^{\text{HF}} = E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] \quad (\text{medium})$$

- **Beyond HF (e.g. RPA/MP2):** correlation is a functional of all MOs  $\psi_n(\mathbf{r})$  and eigenvalues  $\varepsilon_n$ :

$$E_{\text{xc}}^{\text{RPA/MP2}} = E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] + E_{\text{c}}^{\text{RPA/MP2}}[\psi_1, \dots, \psi_{N_{\text{occ}}}, \psi_{N_{\text{occ}}+1}, \dots] \quad (\text{expensive})$$

RPA and MP2 total energy calculations are postprocessing procedures:

- 1 Converge an SCF (for MP2 typically pure HF, for RPA typically PBE):

$$\left( -\frac{\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{Hartree}}(\mathbf{r}) + v_{\text{xc}}(\mathbf{r}) \right) \psi_n(\mathbf{r}) = \varepsilon_n \psi_n(\mathbf{r})$$

- 2 Take occupied and virtual molecular orbitals  $\psi_n(\mathbf{r})$  and their energies  $\varepsilon_n$  of converged SCF and calculate total energy:

$$\text{MP2: } E_{\text{tot}} = E_{\text{kin, el-core, Hartree}} + E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] + E_{\text{c}}^{\text{MP2}}[\text{all } \psi_n]$$

$$\text{RPA: } E_{\text{tot}} = E_{\text{kin, el-core, Hartree}} + E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] + E_{\text{c}}^{\text{RPA}}[\text{all } \psi_n]$$

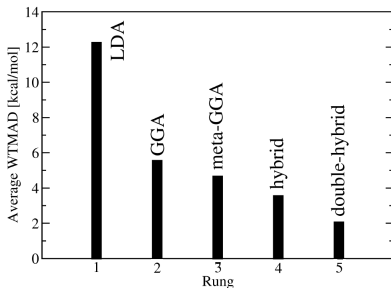
# Motivation: Jacob's Ladder



Jacob's ladder of density functional approximations

## Model of exchange-correlation term improved over time

- Pure DFT is not enough
- Hybrid schemes including HF improve
- Chemical accuracy (1 kcal/mol) not reached
- Go beyond DFT and hybrids with double hybrids containing MP2 and RPA



Accuracy of the density functional approximations

Mundy *et al.* SCIDAC 2010

L. Goerigk, S. Grimme, Phys. Chem. Chem. Phys. **13**, 6670 (2011)

- For MP2 and RPA, converge an SCF with DFT or hybrid to get  $\psi_n, \epsilon_n$ .
- Both methods, MP2 and RPA do not give the exact correlation energy. RPA and MP2 are two different approximations to the correct correlation energy.
- Second order Møller-Plesset perturbation theory (MP2):

$$E_c^{\text{MP2}} = - \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \frac{(ia|jb)[2(ia|jb) - (ib|ja)]}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j}$$

- Random-phase approximation (RPA):

$$E_c^{\text{RPA}} = - \frac{1}{2\pi} \int_0^\infty d\omega \text{Tr}[\ln(1 - \chi(i\omega)v) + \chi(i\omega)v]$$

$$\chi(\mathbf{r}_1, \mathbf{r}_2, i\omega) = 2 \sum_i^{\text{occ}} \sum_a^{\text{virt}} \psi_i(\mathbf{r}_1)\psi_a(\mathbf{r}_1)\psi_i(\mathbf{r}_2)\psi_a(\mathbf{r}_2) \frac{\epsilon_i - \epsilon_a}{(\epsilon_i - \epsilon_a)^2 + \omega^2}, \quad v(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

- In both cases 4-center 2-electron repulsion integrals  $(ia|jb)$  have to be computed

$$(ia|jb) = \iint \psi_i(\mathbf{r}_1)\psi_a(\mathbf{r}_1) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \psi_j(\mathbf{r}_2)\psi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

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Gaussian Auxiliary Basis  $\{P\}$ ; Resolution of the Identity (RI):

$$(ia|jb) \approx (ia|jb)_{RI} = \sum_{PQ} (ia|P) V_{PQ}^{-1} (Q|jb) = \sum_Q B_{ia}^Q B_{jb}^Q, \quad B_{ia}^Q = \sum_P (ia|P) V_{PQ}^{-1/2}$$

where

$$(ia|P) = \iint \psi_i(\mathbf{r}_1) \psi_a(\mathbf{r}_1) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \varphi_P(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

$$V_{PQ} = \iint \varphi_P(\mathbf{r}_1) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \varphi_Q(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Instead of computing 4-center 2-electron repulsion integrals  $(ia|jb)$ , we compute 3-center 2-electron repulsion integrals!

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$$E_C^{\text{MP2}} = - \sum_{ij}^{\text{occ}} \sum_{ab}^{\text{virt}} \frac{(ia|jb) [2(ia|jb) - (ib|ja)]}{\varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j}$$

- Calculate Integrals using RI

$$(ia|jb) = \sum_P B_{ia}^P B_{jb}^P \quad (1)$$

- Scaling of computational effort is  $\mathcal{O}(N^5)$ , see Eq. (1):

occ. orbitals  $(i, j)$  each  $\sim N$

virt. orbitals  $(a, b)$  each  $\sim N$

auxiliary basis  $(P) \sim N$

$$E_c^{\text{MP2,OS}} = - \sum_{iajb} \frac{(ia|jb)^2}{\varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j} = - \sum_{iajb} \frac{(ia|jb)^2}{\Delta_{ij}^{ab}}$$

- Empirical observation [Y. Jung *et al.*, J. Chem. Phys. **121**, 9793 (2004)]:

$$E_c^{\text{MP2}} \approx 1.3 \cdot E_c^{\text{MP2,OS}}$$

- Transform  $1/x = \int_0^\infty e^{-xt} dt$  (I) and numerical integration (II)

$$E_c^{\text{MP2,OS}} \stackrel{\text{(I)}}{=} - \int_0^\infty dt \sum_{iajb} (ia|jb)^2 e^{-t\Delta_{ij}^{ab}} \stackrel{\text{(II)}}{\approx} - \sum_q^{N_q} w_q \sum_{iajb} \underbrace{(ia|jb)^2 e^{-t_q \Delta_{ij}^{ab}}}_{=: A_q}$$

- With RI the overall scaling is reduced to  $\mathcal{O}(N^4)$  (cheaper than RI-MP2!)

$$A_q = \sum_{PQiajb} B_{ia}^P B_{jb}^P B_{ia}^Q B_{jb}^Q e^{-t_q \Delta_{ij}^{ab}} = \sum_{PQ} \sum_{ia} B_{ia}^P B_{ia}^Q e^{-t_q(\varepsilon_a - \varepsilon_i)} \sum_{jb} B_{jb}^P B_{jb}^Q e^{-t_q(\varepsilon_b - \varepsilon_j)}$$

- $N_q = 6-8$  quadrature points  $t_q$  give  $\mu\text{H}$  accuracy

RPA correlation energy in matrix representation:

$$E_c^{\text{RPA}} = \int_0^\infty \frac{d\omega}{2\pi} \text{Tr}(\ln(\mathbf{1} + \mathbf{Q}(\omega)) - \mathbf{Q}(\omega))$$

where  $\mathbf{Q}(\omega)$  is a matrix of size  $N_{\text{RI}} \times N_{\text{RI}}$  with entries

$$Q_{PQ}(\omega) = 2 \sum_{ia} B_{ia}^P \frac{\varepsilon_a - \varepsilon_i}{(\varepsilon_a - \varepsilon_i)^2 + \omega^2} B_{ia}^Q.$$

- RI-RPA overall scaling  $\mathcal{O}(N^4)$
- 10-15 quadrature points (minimax quadrature) for  $\omega$  give  $\mu\text{H}$  accuracy
- very recently:  $\mathcal{O}(N^3)$  scaling RPA with some tricks

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- On 64 cores CRAY XC30:

System	Atoms	Basis (RI basis)	MP2	SOS-MP2	RPA
64 water molec.	192	3648 (8704)	140 min	64 min	48 min

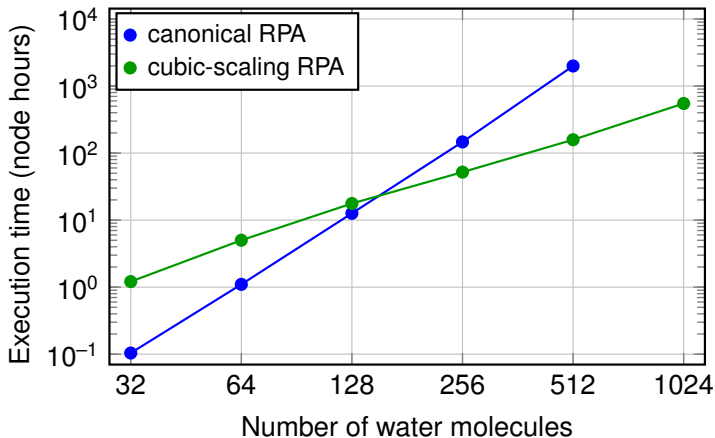
Timings on 64 cores CRAY CX30 (= 8 nodes on Piz Daint, smallest possible number of nodes due to memory requirements), TZVP basis set, quadrature points: 16 (RPA), 10 (OS-MP2).

- On 32768 cores CRAY XC30:

System	Atoms	Basis (RI basis)	RPA [ $\mathcal{O}(N^4)$ -scaling]
64 water molecules	192	3648 (8704)	22 s
512 water molecules	1536	29184 (69632)	1.7 h

Timings on 32768 cores CRAY CX30 (= 4096 nodes on Piz Daint, 78 % of Piz Daint, worldwide 6th largest supercomputer), TZVP basis set, quadrature points: 16.

# Scaling of $\mathcal{O}(N^4)$ (state of the art) and $\mathcal{O}(N^3)$ RPA



Scaling of quartic-scaling RPA (blue) and cubic-scaling RPA (green) with the system size, basis: cc-TZVP, CRAY XC40 machine with 36 cores per node.



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# Input: Basis set and RI basis set for RI-MP2 and RI-RPA

```
&DFT
  BASIS_SET_FILE_NAME ../BASIS_RI_cc-TZ      ! RI basis sets are available for the cc-TZVP basis
  POTENTIAL_FILE_NAME ../POTENTIAL

      .
      .
      .

&END DFT
&SUBSYS

      .
      .
      .

&KIND H
  BASIS_SET cc-TZ
  RI_AUX_BASIS_SET RI_TZ      ! For RI calculations, we additionally need RI basis sets.
  POTENTIAL GTH-PBE-q1
&END KIND
&KIND O
  BASIS_SET cc-TZ
  RI_AUX_BASIS_SET RI_TZ      ! For RI calculations, we additionally need RI basis sets.
  POTENTIAL GTH-PBE-q6
&END KIND
&END SUBSYS
```

# Input non-periodic\* MP2 on top of Hartree-Fock-SCF

\*for periodic calculations, change section PERIODIC NONE and truncate Hartree-Fock

```
&DFT
BASIS_SET_FILE_NAME ../BASIS_RI_cc-TZ
POTENTIAL_FILE_NAME ../POTENTIAL
&QS
METHOD GPW
&END QS
&POISSON
PERIODIC NONE
PSOLVER MT
&END POISSON
&SCF
SCF_GUESS RESTART
&END SCF
&XC
&XC_FUNCTIONAL NONE
&END XC_FUNCTIONAL
&HF
FRACTION 1.00
&SCREENING
EPS_SCHWARZ 1.0E-8
&END SCREENING
&END HF
&END XC
&WF_CORRELATION
METHOD RI_MP2_GPW ! do MP2 calculation
NUMBER_PROC 2 ! number of processes for storing grids. In case you run out of memory when
! computing integrals, increase number. Large number slows the calculation down.
&RI_MP2
&END RI_MP2
&END WF_CORRELATION
&END DFT
```

# Input non-periodic SOS-MP2 on top of Hartree-Fock-SCF

```
&DFT
BASIS_SET_FILE_NAME ../BASIS_RI_cc-TZ
POTENTIAL_FILE_NAME ../POTENTIAL
&QS
METHOD GPW
&END QS
&POISSON
PERIODIC NONE
PSOLVER MT
&END POISSON
&SCF
SCF_GUESS RESTART
&END SCF
&XC
&XC_FUNCTIONAL NON
&END XC_FUNCTIONAL
&HF
FRACTION 1.00
&SCREENING
EPS_SCHWARZ 1.0E-8
&END SCREENING
&END HF
&END XC
&WF_CORRELATION
METHOD RI_SOS_LAPLACE ! do SOS-MP2 calculation
NUMBER_PROC 2
SCALE_S 1.3 ! scale the OS-MP2 energy by 1.3
&RI_LAPLACE
QUADRATURE_POINTS 10
SIZE_INTEG_GROUP 64 ! number of processes dealing with a single frequency grid point. Increase
! this number when running out of memory during the frequency integration.
! A larger number slows the calculation down.
&END RI_LAPLACE
&END WF_CORRELATION
&END DFT
```

# Input non-periodic RPA on top of PBE-SCF

```
&DFT
BASIS_SET_FILE_NAME ../BASIS_RI_cc-TZ
POTENTIAL_FILE_NAME ../POTENTIAL
&QS
METHOD GPW
&END QS
&POISSON
PERIODIC NONE
PSOLVER MT
&END POISSON
&SCF
SCF_GUESS RESTART
&END SCF
&XC
&XC_FUNCTIONAL PBE
&END XC_FUNCTIONAL
&END XC
&WF_CORRELATION
METHOD RI_RPA_GPW ! do an RPA calculation
NUMBER_PROC 2
&RI_RPA
MINIMAX_QUADRATURE TRUE ! choose Minimax quadrature grid points. Then, only few grid points are
! necessary in order to converge the frequency integration in Eq. (2).
QUADRATURE_POINTS 15 ! number of grid points for the frequency integration
SIZE_FREQ_INTEG_GROUP 16 ! number of processes dealing with a single frequency grid point. Increase
! this number when running out of memory during the frequency integration.
! A larger number slows the calculation down.
:
:
:
```

# Optimize RI basis (in case you want other basis than cc-TZVP)

```
&XC
&XC_FUNCTIONAL
  &PBE
    SCALE_X 0.00
    SCALE_C 0.00
  &END
&END XC_FUNCTIONAL
&HF
  FRACTION 1.00
  &SCREENING
    EPS_SCHWARZ 1.0E-8
    SCREEN_ON_INITIAL_P FALSE
  &END
&WF_CORRELATION
  METHOD OPTIMIZE_RI_BASIS
  &OPT_RI_BASIS
    DELTA_I_REL 5.0E-6
    DELTA_RI 5.0E-5
    EPS_DERIV 1.0E-3
    MAX_ITER 100
  BASIS_SIZE MEDIUM ! in case you do not want to specify the number of auxiliary basis functions
  NUM_FUNC 4 4 3 2 1 ! specify number of basis functions in RI basis for each shell, e.g.
                    ! there are 3 d-functions and 1 g-function. Rule of thumb: maximum
                    ! 1-quantum number of RI basis exceeds the primary basis by one.
&END
  NUMBER_PROC 1
&END
&END XC

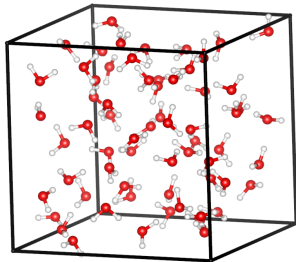
.
.
.
```

```
&END DFT
&SUBSYS

.
.
.
&KIND H
  BASIS_SET cc-QZVP ! basis set for which
                    ! RI basis is optimized
  POTENTIAL GTH-PBE-q1
&END KIND
&COORD ! optimize RI basis for H2 molecule
        ! and take it for arbitrary environ-
        ! ment of H
  H 0.0000 0.0000 0.0000
  H 0.0000 0.0000 1.42000
&END COORD
&END SUBSYS
```

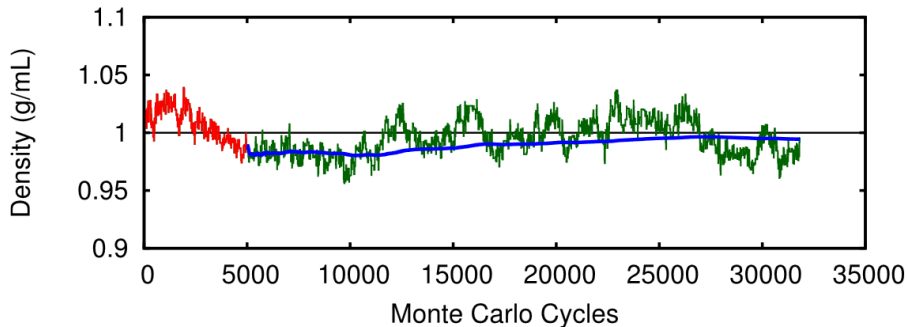
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- **Goal:** Determine **density** of liquid water at ambient conditions.
- **Method:** Isobaric–Isothermal Monte Carlo  
Pre-sampling using fast methods (classical force fields, DFT with PBE)
- 64 water molecules, 192 atoms, 256 active electrons



- cc-TZVP Basis, [3s3p2d1f], [3s2p1d], 3648 basis functions, 8704 RI basis functions





- NpT simulations ( $T = 295$  K and  $p = 1$  bar)

[M. Del Ben, J. Hutter, and J. VandeVondele: Probing the structural and dynamical properties of liquid water with models including non-local electron correlation, *J. Chem. Phys.* **143**, 054506 (2015)]

	$\rho$ [g/mL]	Error estimate $\Delta_\rho$
BLYP	0.797	0.008
BLYP-D3	1.066	0.007
PBE-D3	1.055	0.006
M062X-D3 (ADMM)	1.004	0.008
MP2	1.020	0.004
RPA	0.994	0.006
Exp.	0.998	

[M. Del Ben, J. Hutter, and J. VandeVondele: Probing the structural and dynamical properties of liquid water with models including non-local electron correlation, J. Chem. Phys. **143**, 054506 (2015)]

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# Quasiparticle energies in GW: Theory

- In DFT and Hartree-Fock, there is no theoretical foundation that the eigenvalues  $\varepsilon_n$  from an SCF,

$$\left( -\frac{\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{H}}(\mathbf{r}) + v_{\text{xc}}(\mathbf{r}) \right) \psi_n(\mathbf{r}) = \varepsilon_n \psi_n(\mathbf{r})$$

have anything to do with single-electron energies.

- Theorem:**

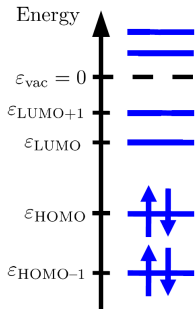
A self-energy  $\Sigma(\mathbf{r}, \mathbf{r}', \varepsilon)$  (non-local, energy-dependent) containing exchange and correlation effects exists, such that the solution of

$$\left( \frac{-\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{H}}(\mathbf{r}) \right) \psi_n(\mathbf{r}) + \int d\mathbf{r}' \Sigma(\mathbf{r}, \mathbf{r}', \varepsilon_n) \psi_n(\mathbf{r}') = \varepsilon_n \psi_n(\mathbf{r})$$

gives the correct quasiparticle energies  $\varepsilon_n$  of the interacting many-electron system.

- In the GW approximation, the self-energy reads

$$\Sigma^{GW}(\mathbf{r}, \mathbf{r}', \varepsilon) = \frac{i}{2\pi} \int_{-\infty}^{\infty} d\varepsilon' G(\mathbf{r}, \mathbf{r}', \varepsilon - \varepsilon', \{\varepsilon_n\}, \{\psi_n\}) W(\mathbf{r}, \mathbf{r}', \varepsilon', \{\varepsilon_n\}, \{\psi_n\})$$



Single-electron (quasiparticle) levels  
of a closed shell molecule

Reminder: One gets the correct eigenvalues  $\varepsilon_n$  by solving

$$\left( \frac{-\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{H}}(\mathbf{r}) \right) \psi_n(\mathbf{r}) + \int d\mathbf{r}' \Sigma(\mathbf{r}, \mathbf{r}', \varepsilon_n) \psi_n(\mathbf{r}') = \varepsilon_n \psi_n(\mathbf{r})$$

- 1 Practically, first converge SCF with DFT

$$\left( -\frac{\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{H}}(\mathbf{r}) + v_{\text{xc}}(\mathbf{r}) \right) \psi_n^{\text{DFT}}(\mathbf{r}) = \varepsilon_n^{\text{DFT}} \psi_n^{\text{DFT}}(\mathbf{r})$$

- 2 Compute the  $(n, n)$  diagonal elements  $\Sigma_n^{\text{GW}}$  of the  $G_0W_0$  self-energy and the xc potential:

$$\Sigma_n^{\text{GW}} = \iint d\mathbf{r} d\mathbf{r}' \psi_n^{\text{DFT}}(\mathbf{r}) \Sigma^{\text{GW}}(\mathbf{r}, \mathbf{r}', \varepsilon_n^{\text{DFT}}) \psi_n^{\text{DFT}}(\mathbf{r}')$$

$$V_{\text{xc},n} = \int d\mathbf{r} \psi_n^{\text{DFT}}(\mathbf{r}) v_{\text{xc}}(\mathbf{r}) \psi_n^{\text{DFT}}(\mathbf{r})$$

- 3 Correct the DFT eigenvalues  $\varepsilon_n^{\text{DFT}}$  by  $G_0W_0$  ( $Z_n \approx 1$  typically):

$$\varepsilon_n^{G_0W_0} = \varepsilon_n^{\text{DFT}} + Z_n \left( \Sigma_n^{\text{GW}} - V_{\text{xc},n} \right)$$

Eigenvalue selfconsistent GW (evGW) yields best results for the HOMO-LUMO gap of molecules [Blase et al., Phys. Rev. B 83, 115103 (2011)]:

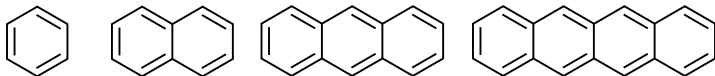
Iterate the  $G_0W_0$  equation to self-consistency in the eigenvalues:

$$G_0W_0: \quad \varepsilon_n^{G_0W_0} = \varepsilon_n^{\text{DFT}} + Z_n \left( \Sigma_n^{\text{GW}}(\{\varepsilon_n^{\text{DFT}}\}, \{\psi_n^{\text{DFT}}\}) - v_{\text{xc},n} \right)$$

$$\text{evGW(1):} \quad \varepsilon_n^{\text{evGW(1)}} = \varepsilon_n^{G_0W_0} + Z_n \left( \Sigma_n^{\text{GW}}(\{\varepsilon_n^{G_0W_0}\}, \{\psi_n^{\text{DFT}}\}) - \Sigma_n^{\text{GW}}(\{\varepsilon_n^{\text{DFT}}\}, \{\psi_n^{\text{DFT}}\}) \right)$$

$$\text{evGW(2):} \quad \varepsilon_n^{\text{evGW(2)}} = \varepsilon_n^{\text{evGW(1)}} + Z_n \left( \underbrace{\Sigma_n^{\text{GW}}(\{\varepsilon_n^{\text{evGW(1)}}\}, \{\psi_n^{\text{DFT}}\}) - \Sigma_n^{\text{GW}}(\{\varepsilon_n^{G_0W_0}\}, \{\psi_n^{\text{DFT}}\})}_{\rightarrow 0 \text{ in case of self-consistency}} \right)$$

```
&DFT
BASIS_SET_FILE_NAME ./BASIS
POTENTIAL_FILE_NAME ./POTENTIAL
UKS
MULTIPLICITY 1
&QS
METHOD GPW
&END QS
&POISSON
PERIODIC NONE
&END POISSON
&XC
&XC_FUNCTIONAL PBE
&END XC_FUNCTIONAL
&WF_CORRELATION
METHOD RI_RPA_GPW
&RI_RPA
RPA_NUM_QUAD_POINTS 100
SIZE_FREQ_INTEG_GROUP 512
GW
&HF
FRACTION 1.0000000
&SCREENING
EPS_SCHWARZ 1.0E-11
SCREEN_ON_INITIAL_P FALSE
&END SCREENING
&MEMORY
MAX_MEMORY 0
&END
&END HF
:
:
&RI_GOW0
! tell the code how many eigenvalues
! should be corrected by GW
CORR_OCC 1000
CORR_VIRT 400
! how many evGW iterations should be
! done at most
EV_SC_ITER 10
&END RI_GOW0
&END RI_RPA
NUMBER_PROC 16
&END
&END XC
&END DFT
:
:
:
```



First four linear acenes.

HOMO-LUMO gaps of known acenes from DFT,  $G_0W_0$  and *evGW* starting from two functionals:

$N_{\text{rings}}$	PBE0			tuned CAM-B3LYP			Exp.
	DFT	$G_0W_0$	<i>evGW</i>	DFT	$G_0W_0$	<i>evGW</i>	
3	3.87	6.63	6.94	6.16	6.90	6.95	6.91
4	3.03	5.54	5.83	5.30	5.81	5.84	5.91
5	2.61	4.77	5.11	4.73	5.01	5.03	5.22



- 1 Beyond-Hartree-Fock methods: Definition, benefits and limitations
- 2 The resolution of the identity (RI)
- 3 RI-MP2, RI-SOS-MP2 and RI-RPA for the correlation energy
- 4 Execution time of RI-MP2, RI-SOS-MP2 and RI-RPA
- 5 Inputs for RI-MP2, RI-SOS-MP2 and RI-RPA
- 6 Application: Density of water from RI-MP2 and RI-RPA Monte Carlo simulations
- 7 RI-GW for single-particle energies: Definition, input and application
- 8 Concluding remarks**

## MP2 and RPA implemented by Mauro Del Ben:

- **MP2**: M. Del Ben, J. Hutter, and J. VandeVondele, JCTC **8**, 4177-4188 (2012).
- **RI-MP2, RI-RPA**: M. Del Ben, J. Hutter, and J. VandeVondele, JCTC **9**, 2654-2671 (2013).
- **More efficient and massively parallel RI-RPA with GPUs**: M. Del Ben, O. Schütt, T. Wentz, P. Messmer, J. Hutter, and J. VandeVondele, Comp. Phys. Comm. **187**, 120-129 (2015).
- **RI-MP2 forces**: M. Del Ben, J. Hutter, and J. VandeVondele, JCP **143**, 102803 (2015).

## Open-shell MP2 forces implemented by Vladimir Rybkin:

- V. V. Rybkin, and J. VandeVondele, JCTC **12**, 2214-2223 (2016).

## GW and cubic-scaling RPA by myself:

- **GW**: J. Wilhelm, M. Del Ben, and J. Hutter, JCTC **12**, 3623-3635 (2016).
- **Cubic-scaling RI-RPA**: J. Wilhelm, P. Seewald, M. Del Ben, and J. Hutter, submitted.

- MP2 and RPA: two different approximate approaches for the correlation energy
  - post-SCF methods: first, converge SCF with Hartree-Fock (for MP2) or method of your choice (for RPA), then calculate total energy with MP2, RPA
  - accurate total energies, e. g. for MD to determine density of water
- *GW*: method to compute single-electron energies
- High cost: MP2:  $\mathcal{O}(N^5)$ , SOS-MP2, RPA, *GW*:  $\mathcal{O}(N^4)$ , RPA also  $\mathcal{O}(N^3)$ 
  - RPA for 64 water molecules on 64 CRAY CX30 cores: 48 min
  - largest RPA calculation so far: 1024 water molecules

## Where you may need MP2, RPA and GW

- MP2 and RPA:
  - 1 Accurate correlation energy needed. Then use double hybrid functionals containing MP2 and RPA.
  - 2 Systems where van-der-Waals interaction is important and dispersion-corrected density functionals do not give accurate results. Maybe pure MP2 or pure RPA give more accurate results.
- GW: For calculating single-electron levels, e.g. the HOMO-LUMO gap (there is no foundation to use DFT to compute single-electron levels!).

### Offer

In case you want to do an MP2, RPA, or GW calculation, you are cordially invited to write me an e-mail. I am happy to help!