Correlated Wave Function Methods in CP2K: Second-Order Møller-Plesset Perturbation Theory (MP2) and Random-Phase Approximation (RPA)

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Beyond-Hartree-Fock methods: Definition, benefits and limitations

- 2 The resolution of the identity (RI)
- 8 RI-MP2, RI-SOS-MP2 and RI-RPA for the correlation energy
- Al-MP2 analytical forces
- Execution time of RI-MP2, RI-SOS-MP2 and RI-RPA
- Inputs for RI-MP2, RI-SOS-MP2 and RI-RPA
- Applications
- Concluding remarks





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The exchange-correlation term in Beyond-Hartree-Fock methods

Total energy *E*_{tot} of electrons in a molecule or periodic system:

$$E_{\text{tot}} = E_{\text{kin}} + E_{\text{el-core}} + E_{\text{Hartree}} + E_{\text{exchange (x)}} + E_{\text{correlation (c)}}$$

• **DFT**: exchange and correlation is a functional of the density $n(\mathbf{r}) = \sum_{i=1}^{N_{occ}} |\psi_i(\mathbf{r})|^2$:

$$E_{\rm xc}^{\rm DFT} = E_{\rm xc}^{\rm DFT}[n]$$
 (cheap)

• Hartree-Fock (HF): No correlation, exchange energy only depends on occupied MOs $\psi_i(\mathbf{r})$:

$$E_{\rm xc}^{\rm HF} = E_{\rm x}^{\rm HF}[\psi_1, \dots, \psi_{N_{\rm occ}}] \tag{medium}$$

Beyond HF (e.g. RPA/MP2): correlation is a functional of all MOs ψ_n(r) and eigenvalues ε_n:

$$E_{\rm xc}^{\rm RPA/MP2} = E_{\rm x}^{\rm HF}[\psi_1, \dots, \psi_{N_{\rm occ}}] + E_{\rm c}^{\rm RPA/MP2}[\psi_1, \dots, \psi_{N_{\rm occ}}, \psi_{N_{\rm occ}+1}, \dots] \qquad (\text{expensive})$$

RPA and MP2 total energy calculations are postprocessing procedures:

Onverge an SCF (for MP2 always pure HF, for RPA typically PBE):

$$\left(-\frac{\nabla^2}{2} + v_{\text{el-core}}(\mathbf{r}) + v_{\text{Hartree}}(\mathbf{r}) + v_{\text{xc}}(\mathbf{r})\right)\psi_n(\mathbf{r}) = \varepsilon_n \psi_n(\mathbf{r})$$

Take occupied and virtual molecular orbitals ψ_n(**r**) and their energies ε_n of converged SCF and calculate total energy:

MP2:
$$E_{\text{tot}} = E_{\text{kin, el-core, Hartree}} + E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] + E_{\text{c}}^{\text{MP2}}[\text{all } \psi_n]$$

RPA: $E_{\text{tot}} = E_{\text{kin, el-core, Hartree}} + E_{\text{x}}^{\text{HF}}[\psi_1, \dots, \psi_{N_{\text{occ}}}] + E_{\text{c}}^{\text{RPA}}[\text{all } \psi_n]$

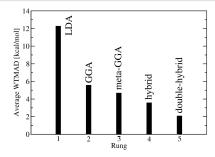


Motivation: Jacob's Ladder



Model of exchange-correlation term improved over time

- Pure DFT is not enough
- Hybrid schemes including HF improve
- Chemical accuracy (1 kcal/mol) not reached
- Go beyond DFT and hybrids with double hybrids containing MP2 and RPA



Accuracy of the density functional approximations

Mundy et al. SCIDAC 2010

L. Goerigk, S. Grimme, Phys. Chem. Chem. Phys. 13, 6670 (2011)



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MP2 and RPA correlation energies

- For MP2 and RPA, converge an SCF with HF/DFT to get ψ_n, ε_n .
- Both methods, MP2 and RPA do not give the exact correlation energy. RPA and MP2 are two
 different approximations to the correct correlation energy.
- Second order Møller-Plesset perturbation theory (MP2):

$$E_{\rm c}^{\rm MP2} = -\sum_{ij}^{\rm occ} \sum_{ab}^{\rm virt} \frac{(ia|jb)[2(ia|jb) - (ib|ja)]}{\varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j}$$

• Random-phase approximation (RPA):

$$E_{
m c}^{
m RPA} = - rac{1}{2\pi} \int_{0}^{\infty} d\omega \; {
m Tr}[{
m ln}(1-\chi(i\omega)v)+\chi(i\omega)v]$$

$$\chi(\mathbf{r}_1, \mathbf{r}_2, i\omega) = 2\sum_{i}^{\text{occ}} \sum_{a}^{\text{virt}} \psi_i(\mathbf{r}_1) \psi_a(\mathbf{r}_1) \psi_i(\mathbf{r}_2) \psi_a(\mathbf{r}_2) \frac{\varepsilon_i - \varepsilon_a}{(\varepsilon_i - \varepsilon_a)^2 + \omega^2}, \quad v(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

• In both cases 4-center 2-electron repulsion integrals (*ia*|*jb*) have to be computed

$$(ia|jb) = \iint \psi_i(\mathbf{r}_1)\psi_a(\mathbf{r}_1)\frac{1}{|\mathbf{r}_1-\mathbf{r}_2|}\psi_j(\mathbf{r}_2)\psi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$



- Rayleigh-Schrödinger perturbation theory:
- Zeroth-order Hamiltonian is the Fock operator: $\hat{H}_0 = \hat{F}$,
- Perturbation is the fluctuation potential: $\hat{V} = \hat{H} \hat{F}$,
- Zeroth-order energy is the sum of orbital energies: $\sum_i \epsilon_i$
- First-order energy: HF energy (NO MP1 EXISTS!)
- Second-order energy: MP2 energy

$$E_{\rm c}^{\rm MP2} = -\sum_{ij}^{\rm occ} \sum_{ab}^{\rm virt} \frac{(ia|jb)t_{ij}^{ab}}{\varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j}$$

- The simplest ab inition correlated wave function method: involves multiple Slater determinants
- Also known as many-body perturbation theory (MBPT), although there are other variants of MBPT (e.g. GW)



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Gaussian Auxiliary Basis {*P*}; Resolution of the Identity (RI), or density fitting: $(ia|jb) \approx (ia|jb)_{RI} = \sum_{PQ} (ia|P) V_{PQ}^{-1}(Q|jb) = \sum_{Q} B_{ia}^{Q} B_{jb}^{Q}, \quad B_{ia}^{Q} = \sum_{P} (ia|P) V_{PQ}^{-1/2}$

where

$$(ia|P) = \iint \psi_i(\mathbf{r}_1)\psi_a(\mathbf{r}_1)\frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|}\varphi_P(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$
$$V_{PQ} = \iint \varphi_P(\mathbf{r}_1)\frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|}\varphi_Q(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Instead of computing 4-center 2-electron repulsion integrals (*ia*|*jb*), we compute 3-center 2-electron repulsion integrals!



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$$E_{c}^{MP2} = -\sum_{ij}^{occ} \sum_{ab}^{virt} \frac{(ia|jb) [2(ia|jb) - (ib|ja)]}{\varepsilon_{a} + \varepsilon_{b} - \varepsilon_{i} - \varepsilon_{j}}$$

• Calculate Integrals using RI

$$(\textit{ia}|\textit{jb}) = \sum_{P} B^{P}_{\textit{ia}} B^{P}_{\textit{jb}}$$

• Scaling of computational effort is $\mathcal{O}(N^5)$, see Eq. (1):

occ. orbitals (i, j) each $\sim N$ virt. orbitals (a, b) each $\sim N$ auxiliary basis $(P) \sim N$



(1)

RI-SOS-MP2: scaled-opposite-spin, Laplace-transformed

$$E_{\rm c}^{\rm MP2,OS} = -\sum_{iajb} \frac{(ia|jb)^2}{\varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j} = -\sum_{iajb} \frac{(ia|jb)^2}{\Delta_{ij}^{ab}}$$

• Empirical observation [Y. Jung et al., J. Chem. Phys. 121, 9793 (2004)]:

$$E_{
m c}^{
m MP2}pprox$$
 1.3 \cdot $E_{
m c}^{
m MP2,OS}$

• Laplace-transform $1/r = \int_0^\infty e^{-rt} dt$ (I) and numerical integration (II)

$$E_{c}^{\text{MP2,OS}} \stackrel{(I)}{=} -\int_{0}^{\infty} dt \sum_{iajb} (ia|jb)^{2} e^{-t\Delta_{ij}^{ab}} \stackrel{(II)}{\approx} -\sum_{q}^{N_{q}} w_{q} \sum_{iajb} \underbrace{(ia|jb)^{2} e^{-t_{q}\Delta_{ij}^{ab}}}_{=:A_{q}}$$

• With RI the overall scaling is reduced to $\mathcal{O}(N^4)$ (cheaper than RI-MP2!)

$$A_{q} = \sum_{PQiajb} B_{ia}^{P} B_{jb}^{Q} B_{ia}^{Q} B_{jb}^{Q} e^{-t_{q} \Delta_{ij}^{ab}} = \sum_{PQ} \sum_{ia} B_{ia}^{P} B_{ia}^{Q} e^{-t_{q}(\varepsilon_{a} - \varepsilon_{i})} \sum_{jb} B_{jb}^{P} B_{jb}^{Q} e^{-t_{q}(\varepsilon_{b} - \varepsilon_{j})}$$

• $N_q = 6-8$ quadrature points t_q give μ H accuracy

RPA correlation energy in matrix representation:

$$E_{
m c}^{
m RPA} = \int_{0}^{\infty} rac{d\omega}{2\pi} \, {
m Tr} \left({
m In} \left({f 1} + {f Q}(\omega)
ight) - {f Q}(\omega)
ight)$$

where $Q(\omega)$ is a matrix of size $N_{\rm RI} \times N_{\rm RI}$ with entries

$$Q_{PQ}(\omega) = 2 \sum_{ia} B_{ia}^{P} rac{arepsilon_{a} - arepsilon_{i}}{(arepsilon_{a} - arepsilon_{i})^{2} + \omega^{2}} B_{ia}^{Q}.$$

- RI-RPA overall scaling $\mathcal{O}(N^4)$
- 10-15 quadrature points (minimax quadrature) for ω give μ H accuracy
- very recently: $\mathcal{O}(N^3)$ scaling RPA with some tricks



Energy is not variational wrt. orbitals and determinant coefficients:

$$E^{(2)^{(x)}} = 4 \sum_{Q}^{AUX} \sum_{\mu\nu}^{AO} \Gamma^{Q}_{\mu\nu} (\mu\nu | Q)^{x} - 2 \sum_{PQ}^{AUX} \Gamma^{PQ} (P | Q)^{x} + 2 \sum_{Pq}^{SO} [P^{(2)}_{\rho q} F^{x}_{\rho q} - W^{(2)}_{\rho q} S^{x}_{\rho q}],$$

where μ , ν are basis functions, p, q are spin orbitals, ^x is a derivative and

$$\sum_{p}^{SO} = \sum_{p^{lpha}}^{SO^{lpha}} + \sum_{p^{eta}}^{SO^{eta}}.$$

Electron (spin) density:

$$\rho^{s}(\mathbf{r}) = \frac{1}{2} \sum_{\rho^{\alpha} q^{\alpha}} P_{\rho^{\alpha} q^{\alpha}} \phi_{\rho^{\alpha}}(\mathbf{r}) \phi_{q^{\alpha}}(\mathbf{r}) + (-) \frac{1}{2} \sum_{\rho^{\beta} q^{\beta}} P_{\rho^{\beta} q^{\beta}} \phi_{\rho^{\beta}}(\mathbf{r}) \phi_{q^{\beta}}(\mathbf{r}).$$

MP2 electron and spin densities are only available after a force calculation!



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• On 64 cores CRAY XC30:

| System | Atoms | Basis (RI basis) | MP2 | SOS-MP2 | RPA |
|-----------------|-------|------------------|---------|---------|--------|
| 64 water molec. | 192 | 3648 (8704) | 140 min | 64 min | 48 min |

Timings on 64 cores CRAY CX30 (= 8 nodes on Piz Daint, smallest possible number of nodes due to memory requirements), TZVP basis set, quadrature points: 16 (RPA), 10 (OS-MP2).

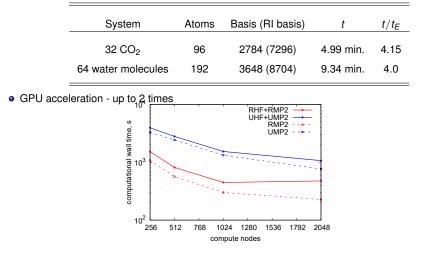
• On 32768 cores CRAY XC30:

| System | Atoms | Basis (RI basis) | RPA [$\mathcal{O}(N^4)$ -scaling] |
|---------------------|-------|------------------|------------------------------------|
| 64 water molecules | 192 | 3648 (8704) | 22 s |
| 512 water molecules | 1536 | 29184 (69632) | 1.7 h |

Timings on 32768 cores CRAY CX30 (= 4096 nodes on Piz Daint, 78% of Piz Daint, worldwide 6th largest supercomputer), TZVP basis set, quadrature points: 16.



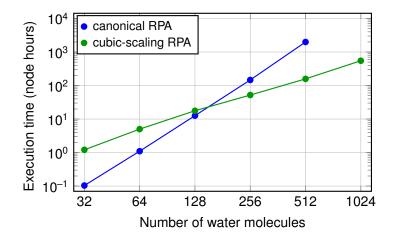
• On 4096 cores CRAX CX 30 with 512 GPUs (=512 nodes on Piz Daint):



Spin-unrestricted (blue) and spin-restricted (red) calculations of 64 water molecules.



Scaling of $\mathcal{O}(N^4)$ (state of the art) and $\mathcal{O}(N^3)$ RPA



Scaling of quartic-scaling RPA (blue) and cubic-scaling RPA (green) with the system size, basis: cc-TZVP, CRAY XC40 machine with 36 cores per node.



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&DFT

BASIS SET FILE NAME ... /BASIS RI CC-TZ ! RI basis sets are available for the cc-TZVP basis POTENTIAL FILE NAME ... /POTENTIAL SEND DFT & SUBSYS SKIND H BASIS SET CC-TZ RI AUX BASIS SET RI TZ ! For RI calculations, we additionally need RI basis sets. POTENTIAL GTH-PBE-q1 SEND KIND &KIND O BASIS_SET cc-TZ RI AUX BASIS SET RI TZ ! For RI calculations, we additionally need RI basis sets. POTENTIAL GTH-PBE-q6 SEND KIND &END SUBSYS



Input non-periodic* MP2 on top of Hartree-Fock-SCF

*for periodic calculations, change section PERIODIC NONE and truncate Hartree-Fock Coulomb operator

```
&DFT
  BASIS SET FILE NAME ... /BASIS RI CC-TZ
 POTENTIAL FILE NAME ... /POTENTIAL
  &OS
   METHOD GPW
  &END OS
  &POISSON
   PERIODIC NONE
   PSOLVER MT
  &END POISSON
  &SCF
    SCF GUESS RESTART
  &END SCF
  exc.
    &XC FUNCTIONAL NONE
    &END XC FUNCTIONAL
    LHE
     FRACTION 1 00
      &SCREENING
       EPS SCHWARZ 1.0E-8
      SCREENING
    SEND HE
  SEND XC
  &WF CORRELATION
   METHOD RI MP2 GPW ! do MP2 calculation
    NUMBER PROC 2
                        ! number of processes for storing grids. In case you run out of memory when
                        ! computing integrals, increase number. Large number slows the calculation down.
   &RI MP2
    &END RI MP2
   &END WF CORRELATION
 SEND DFT
```

University of

Zurich

Input non-periodic SOS-MP2 on top of Hartree-Fock-SCF

```
&DFT
  BASIS SET FILE NAME ... / BASIS RI cc-TZ
  POTENTIAL FILE NAME ... / POTENTIAL
  &QS
   METHOD GPW
  &END OS
  &POISSON
   PERTODIC NONE
   PSOLVER MT
  &END POISSON
  ASCE
    SCF GUESS RESTART
  SEND SCE
  &XC
    &XC FUNCTIONAL NON
    &END XC FUNCTIONAL
    LHE
     FRACTION 1 00
      &SCREENING
        EPS SCHWARZ 1.0E-8
      &END SCREENING
    SEND HE
  SEND XC
  SWF CORRELATION
   METHOD RI SOS LAPLACE ! do SOS-MP2 calculation
   NUMBER PROC 2
   SCALE S 1.3
                   ! scale the OS-MP2 energy by 1.3
   &RI LAPLACE
     QUADRATURE POINTS 10
     SIZE INTEG GROUP 64 ! number of processes dealing with a single frequency grid point. Increase
                            ! this number when running out of memory during the frequency integration.
                            ! A larger number slows the calculation down.
    &END RI LAPLACE
  &END WF CORRELATION
                                                                                                               University of
SEND DFT
                                                                                                               Zurich
```

```
SWF CORRELATION
     METHOD RI MP2 GPW
     &RI MP2
       BLOCK SIZE 1
       EPS CANONICAL 1.0E-7
       FREE HFX BUFFER .FALSE. ! Memory for HFX exchange
     £END
     &CPHF ! Coupled perturbed equations
       EPS CONV 1.0E-6
       MAX ITER 50
     £END
     &WFC GPW
       CUTOFF
                   300
       REL CUTOFF 50
       EPS FILTER 1.0E-12
       EPS GRID 1.0E-10
     £END
     MEMORY 8000
     NUMBER PROC 1
   SEND.
```

- Coupled-perturbed equations for occupied-virtual one-particle density matrix
- Free HF exchange buffer used to solve them: saves time or memory
- Play with memory in the HFX and WF CORRELATION sections



Input non-periodic RPA on top of PBE-SCF

&DFT BASIS SET FILE NAME ... / BASIS RI cc-TZ POTENTIAL FILE NAME ... / POTENTIAL & HF ! Hartree-Fock for exchange energy FRACTION 1 00 &QS METHOD GPW & SCREENING &END QS EPS SCHWARZ 1.0E-8 &END SCREENING &POISSON PERIODIC NONE **EMEMORY** MAX MEMORY 0 ! we need not to store PSOLVER MT FEND POISSON ! any HF integrals because there & SCF ! are no subsequent SCF steps SCF GUESS RESTART SEND MEMORY SEND SCE SEND HE exc. &END RI RPA &XC FUNCTIONAL PBE &END WF CORRELATION SEND XC FUNCTIONAL SEND DET SEND XC **&WF CORRELATION** METHOD RI RPA GPW ! do an RPA calculation NUMBER PROC 2 &RI RPA MINIMAX OUADRATURE TRUE ! choose Minimax quadrature grid points. Then, only few grid points are ! necessary in order to converge the frequency integration in Eq. (2). OUADRATURE POINTS 15 ! number of grid points for the frequency integration SIZE FREO INTEG GROUP 16 ! number of processes dealing with a single frequency grid point. Increase ! this number when running out of memory during the frequency integration. ! A larger number slows the calculation down.



```
&XC
                                                        SEND DFT
  &XC FUNCTIONAL
                                                        & SUBSYS
    &PBE
     SCALE X 0.00
      SCALE C 0.00
    SEND.
  &END XC FUNCTIONAL
                                                         SKIND H
                                                           BASIS SET cc-OZVP
                                                                              ! basis set for which
  SHF
                                                                              ! RI basis is optimized
    FRACTION 1.00
                                                           POTENTIAL GTH-PBE-g1
    SCREENING
                                                         &END KIND
      EPS SCHWARZ 1.0E-8
     SCREEN ON INITIAL P FALSE
                                                         &COORD ! optimize RI basis for H2 molecule
    SEND
                                                                 ! and take it for arbitrary environ-
  SEND
                                                                 1 ment of H
  SWF CORRELATION
                                                           н 0 0000 0 0000
                                                                                0.0000
   METHOD OPTIMIZE RI BASIS
                                                           н 0 0000
                                                                       0 0000 1 42000
    &OPT RI BASIS
                                                         SEND COORD
     DELTA I REL 5.0E-6
                                                        &END SUBSYS
     DELTA RI
                  5 OE-5
     EPS DERIV
                  1.0E-3
     MAX ITER
                  100
      BASIS SIZE MEDIUM ! in case you do not want to specify the number of auxiliary basis functions
     NUM FUNC 4 4 3 2 1 ! specify number of basis functions in RI basis for each shell, e.g.
                         ! there are 3 d-functions and 1 g-function. Rule of thumb: maximum
                         ! 1-guantum number of RI basis exceeds the primary basis by one.
    &END
    NUMBER PROC 1
  $END
SEND XC
```



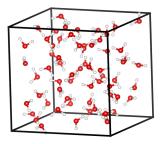
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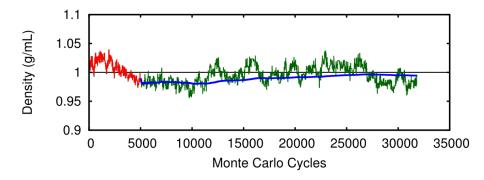
Isobaric–Isothermal Monte Carlo Simulation of Liquid Water

- Goal: Determine density of liquid water at ambient conditions.
- Method: Isobaric–Isothermal Monte Carlo
 Pre-sampling using fast methods (classical force fields, DFT with PBE)
- 64 water molecules, 192 atoms, 256 active electrons



 cc-TZVP Basis, [3s3p2d1f], [3s2p1d], 3648 basis functions, 8704 RI basis functions





• NpT simulations (T = 295 K and p = 1 bar)

[M. Del Ben, J. Hutter, and J. VandeVondele: Probing the structural and dynamical properties of liquid water with models including non-local electron correlation, J. Chem. Phys. **143**, 054506 (2015)]



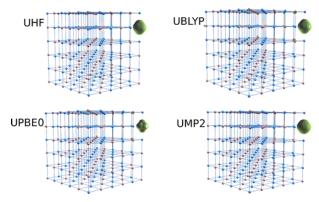
| | ho[g/mL] | Error estimate $\Delta_{ ho}$ |
|-----------------|----------|-------------------------------|
| BLYP | 0.797 | 0.008 |
| BLYP-D3 | 1.066 | 0.007 |
| PBE-D3 | 1.055 | 0.006 |
| M062X-D3 (ADMM) | 1.004 | 0.008 |
| MP2 | 1.020 | 0.004 |
| RPA | 0.994 | 0.006 |
| Exp. | 0.998 | |
| | | |

[M. Del Ben, J. Hutter, and J. VandeVondele: Probing the structural and dynamical properties of liquid water with models including non-local electron correlation, J. Chem. Phys. 143, 054506 (2015)]



F-center in lithium fluoride

- $\bullet~3\times3\times3$ supercell, $Li_{108}F_{107}$
- 10 hours for energy and gradient on 2048 Piz Daint nodes
- Spin-density distributions: MP2 is free of delocalization error present in DFT





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MP2 and RPA implemented by Mauro Del Ben:

- MP2: M. Del Ben, J. Hutter, and J. VandeVondele, JCTC 8, 4177-4188 (2012).
- RI-MP2, RI-RPA: M. Del Ben, J. Hutter, and J. VandeVondele, JCTC 9, 2654-2671 (2013).
- More efficient and massively parallel RI-RPA with GPUs: M. Del Ben, O. Schütt, T. Wentz, P. Messmer, J. Hutter, and J. VandeVondele, Comp. Phys. Comm. 187, 120-129 (2015).
- RI-MP2 forces: M. Del Ben, J. Hutter, and J. VandeVondele, JCP 143, 102803 (2015).

Open-shell MP2 forces implemented by Vladimir Rybkin:

• V. V. Rybkin, and J. VandeVondele, JCTC 12, 2214-2223 (2016).

Cubic-scaling RPA by Jan Wilhelm

• Cubic-scaling RI-RPA: J. Wilhelm, P. Seewald, M. Del Ben, and J. Hutter, submitted.



- MP2 and RPA: two different approximate approaches for the correlation energy
 - post-SCF methods: first, converge SCF with Hartree-Fock (for MP2) or method of your choice (for RPA), then calculate total energy with MP2, RPA
 - accurate total energies, e.g. for MD to determine density of water
- High cost: MP2: $\mathcal{O}(N^5)$, SOS-MP2, RPA: $\mathcal{O}(N^4)$, RPA also $\mathcal{O}(N^3)$
 - RPA for 64 water molecules on 64 CRAY CX30 cores: 48 min
 - largest RPA calculation so far: 1024 water molecules
- Functionality:
 - Energies: MP2, RI-MP2, SOS-MP2, RPA (O(N⁴)), RPA (O(N³))
 - Forces, stress, electron and spin densities: RI-MP2



- Accurate correlation energy needed. Then use double hybrid functionals containing MP2 and RPA
- Systems where van-der-Waals interaction is important and disperion-corrected density functionals do not give accurate results. Maybe pure MP2 or pure RPA give more accurate results
- OFT suffers from delocalization error, e.g. for radicals and electron transfer
- Not for strong/static correlation and bond breaking

